

# Multicompartment micelles formed from star-dendritic triblock copolymers in selective solvents: A dissipative particle dynamics study

Dahuan Liu\*, Chongli Zhong

*Key Lab for Nanomaterials, Ministry of Education, Department of Chemical Engineering, Beijing University of Chemical Technology, 15 BeiSanhuan East Road, Beijing 100029, China*

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## Abstract

Dissipative particle dynamics method was used to study the multicompartment micelles formed from star-dendritic triblock copolymers in selective solvents, in which particular attention was paid to the effects of dendritic structure. The simulations show that the dendritic structure not only influences the morphology and the formation process of multicompartment micelles formed, but also the response of the micelle structure to solvent quality. The information obtained is useful for the future design of multicompartment micelles for practical applications, especially in the field of drug delivery.

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## 1. Introduction

Amphiphilic block copolymer can self-assemble into micelles and vesicles in selective solvent. For the past several decades, these materials with nanometer-scale microphase-separated structures have been applied to many fields, particularly to drug delivery [1,2]. Recently, multicompartment micelle has been introduced as a new class of micelles inspired originally by biological systems, showing great potential applications in biotechnology, multifunctional nanoreactors, medicine, etc. [3,4].

Multicompartment micelles have complex structures consisting of a hydrophilic shell and a hydrophobic core with multidomains, which may extend the applications of traditional micelles to a wide scope. Therefore, after the multicompartment micelles were synthesized successfully, special attentions have been paid to study their morphologies and

structures [4]. As a result, several intriguing morphologies were observed, such as “worm” [5], “hamburger” [6], “sphere-on-sphere” [7], cylinder [8,9], polygonal bilayer sheets [10], and Y-junctions [11]. These works indicated that further investigations are needed to study the influence of the architecture of building block copolymers on the inner structure, as well as towards understanding multicompartment micelle thoroughly [12].

It is worth noting that to date almost all of the multicompartment micelles reported were formed from the linear block copolymers or miktoarm star triblock copolymer. At the same time, another kind of complex polymers, dendrimers [13–17] and hyperbranched copolymers [18], has received increasing attention due to their unique architecture and ability to form self-assembled fascinating structures. Recently, Yan and co-workers [19] have firstly investigated multicompartment micelles from amphiphilic hyperbranched star block copolymers. Unfortunately, the effects of the architecture of dendrimer on the morphology and formation process of the multicompartment micelles are not studied systematically as far as we know and the knowledge on the response to the

\* Corresponding author. Fax: +86 10 64419862.

*E-mail address:* [liudahuan2006@yahoo.com.cn](mailto:liudahuan2006@yahoo.com.cn) (D. Liu).

change of circumstance is still quite lacking, which are pivotal to the practical applications of multicompartment micelles.

Considering that the synthesis of dendrimer as well as the control of architecture of block copolymers systematically are difficult and time-consuming in experiment, computer simulation is a suitable and powerful tool for investigating the behavior of micelles, as well as complementing experimental works in a wider parameter space [20,21]. Therefore, based on the results of star triblock copolymers got in our research group [22], the multicompartment micelles formed from star-dendritic triblock copolymers in selective solvent were studied in this work by a mesoscopic simulation technique, dissipative particle dynamics (DPD) method, which has been successfully used to study the self-assembly behavior of star or linear block copolymers in previous work [21–27]. The main purpose of this work is to try to provide preliminary numerical results and some useful information for further experimental study.

## 2. Method and simulation details

### 2.1. Dissipative particle dynamics method

Dissipative particle dynamics (DPD) method is a mesoscopic simulation technique based on the fluctuation–dissipation theorem [28–30]. In this method, the force acting on a particle contains three pairwise additive parts:

$$\mathbf{f}_i = \sum_{j \neq i} (\mathbf{F}_{ij}^C + \mathbf{F}_{ij}^D + \mathbf{F}_{ij}^R) \quad (1)$$

where  $\mathbf{F}^C$ ,  $\mathbf{F}^D$  and  $\mathbf{F}^R$  represent conservative, dissipative, and random forces, respectively, and the sum runs over all other particles within a certain cutoff radius  $r_c$ . As this is the only length-scale in the system, cutoff radius is always used as the unit of length,  $r_c = 1$ .

The conservative force is given by

$$\mathbf{F}_{ij}^C = \begin{cases} a_{ij}(1 - r_{ij}/r_c)\hat{\mathbf{r}}_{ij} & (r_{ij} < r_c) \\ 0 & (r_{ij} \geq r_c) \end{cases} \quad (2)$$

where  $a_{ij}$  is the maximum repulsion between particles  $i$  and  $j$ , and  $\mathbf{r}_{ij} = \mathbf{r}_i - \mathbf{r}_j$ ,  $r_{ij} = |\mathbf{r}_{ij}|$ , and  $\hat{\mathbf{r}}_{ij} = \mathbf{r}_{ij}/|\mathbf{r}_{ij}|$ . This kind of force is soft repulsion acting along the line of centers between particles.

The dissipative force and the random force are given by

$$\mathbf{F}_{ij}^D = -\gamma\omega^D(r_{ij})(\hat{\mathbf{r}}_{ij} \cdot \mathbf{v}_{ij})\hat{\mathbf{r}}_{ij} \quad (3)$$

$$\mathbf{F}_{ij}^R = \sigma\omega^R(r_{ij})\theta_{ij}\hat{\mathbf{r}}_{ij} \quad (4)$$

where  $\mathbf{v}_{ij} = \mathbf{v}_i - \mathbf{v}_j$ ,  $\omega^D$  and  $\omega^R$  are weight functions vanishing for  $r > r_c$ ,  $\gamma$  is the friction coefficient,  $\sigma$  is the noise amplitude, and  $\theta_{ij}$  is a randomly fluctuating variable with Gaussian statistics.

The two weight functions can be taken simply as

$$\omega^D(r) = [\omega^R(r)]^2 = \begin{cases} (1 - r/r_c)^2 & (r_{ij} < r_c) \\ 0 & (r_{ij} \geq r_c) \end{cases} \quad (5)$$

$$\sigma^2 = 2\gamma k_B T \quad (6)$$

In this work, the relationship between  $a_{ij}$  and the Flory–Huggins  $\chi$ -parameter proposed by Groot and Warren [30] is adopted to determine the repulsion parameter  $a_{ij}$  and then get the conservative force  $F^C$ :

$$a_{ij} = \begin{cases} a_{ii} + 3.27\chi_{ij} & \rho = 3 \\ a_{ii} + 1.45\chi_{ij} & \rho = 5 \end{cases} \quad (7)$$

where  $\rho$  is the density,  $a_{ii}$  is the repulsion parameter between particles of the same type, and its value is derived from the compressibility of pure component by

$$a_{ii} = 75k_B/\rho \quad (8)$$

### 2.2. Simulation details

Multicompartment micelles from the self-assembly of ABC star triblock copolymer have been investigated both in experiment [5] and in simulation work [22,25], and this miktoarm star architecture is believed to be successful remarkably in promoting the formation of multicompartment cores [5]. Therefore, in this work we introduced the dendritic structure into ABC star triblock copolymer to perform preliminary numerical discussion on the multicompartment micelle formed from dendrimers. The numbers of DPD beads associated with A, B, and C blocks are set to 4, 10, and 2 respectively, to ensure the formation of segmented wormlike micelles [22] which have an intriguing structure for loading two or more kinds of agents simultaneously and may be a promising material for drug delivery. As depicted in Fig. 1, model I is in fact the star triblock copolymer to bring into comparison. For model II, the hydrophilic block B is dendritic, while the long hydrophobic block A is dendritic in model III for convenience.

In the previous work [22], we studied the self-assembly of ABC star triblock copolymers into multicompartment micelles in water, and the repulsive interaction parameter  $a_{ij}$  used for the four unlike species was determined by reproducing the experimental observations of Lodge and coworkers [5]:

$$a_{ij} = \begin{pmatrix} & \text{A} & \text{B} & \text{C} & \text{S} \\ \text{A} & 25 & 45 & 75 & 50 \\ \text{B} & 45 & 25 & 90 & 27 \\ \text{C} & 75 & 90 & 25 & 120 \\ \text{S} & 50 & 27 & 120 & 25 \end{pmatrix} \quad (9)$$

where A, B, C and S represent the weakly hydrophobic block, the hydrophilic block, the strongly hydrophobic block, and the solvent, respectively. In this work, DPD repulsion parameters were based on these values.

The DPD simulations were performed in a cubic cell of size  $30 \times 30 \times 30r_c^3$  containing 81,000 DPD beads and periodic boundary conditions were applied. Bead density was taken

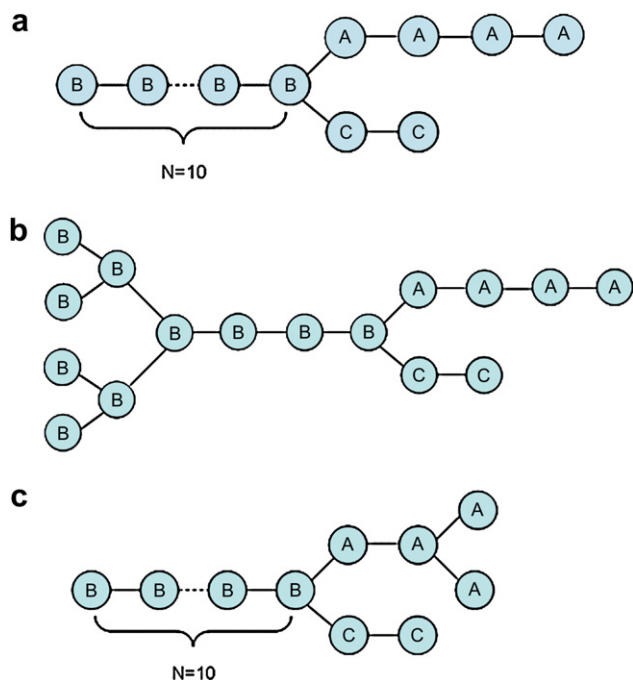


Fig. 1. Schematic structures of star-dendritic block copolymers: (a) model I, (b) model II, and (c) model III.

as 3.0 and the volume fraction of the copolymer was set to be 0.1 to ensure enough micelles can be formed in these systems. For convenience, the cutoff radius  $r_c$ , the particle mass  $m$ , and  $k_B T$  were all taken as unity. The time step  $\Delta t$  was taken as 0.05, and adjacent particles in the polymer chain interacted via a linear spring with a harmonic spring constant of 4.0, and a DPD thermostat with the friction/noise parameter  $\gamma = 4.5$  was used. Depending on the stability of properties of the systems including the mean-square radius of gyration and conservative energy, a total of  $3\text{--}5 \times 10^5$  DPD steps were carried out for each DPD simulation to ensure the equilibration.

### 3. Results and discussion

#### 3.1. Effects of the dendritic structure on the formation of multicompartment micelles

First, the effects of dendritic structure on the formation of multicompartment micelles were studied by comparing the

morphologies of micelles formed from models I–III on the same conditions, as given in Fig. 2.

The simulations show that these three kinds of star-dendritic triblock copolymers can all self-assemble into wormlike multicompartment micelles. Since the length of hydrophobic block A is not long enough, the number of branch points in dendritic part is small in model III, which induces that the difference between the micelles formed from models III and I is negligible. However, the copolymer with hydrophilic dendritic block (model II) tends to form longer wormlike micelle than the other models. Thus, we mainly focus on the cases of models I and II in the following section.

To further understand the effect of dendritic structure, the evolution processes of the corresponding micelle are investigated. Similar to the formation mechanism found in the previous works [22,26], these micelles are formed through fusion process and the general sequence is discrete “hamburger” micelle, short wormlike micelle, and long wormlike micelle. However, from the snapshots of evolution process given in Fig. 3, it can be seen that short wormlike structure and long wormlike structure first appear at 25,000 step (Fig. 3f) and at 87,000 step (Fig. 3h), respectively, in model II, while at 10,000 step (Fig. 3b) and at 35,000 step (Fig. 3d) in model I, that is, the dendritic hydrophilic block leads to the slower fusion process.

Interestingly, in spite of dendritic conformation forming a very protective layer and shielding the hydrophobic core to lead to a slow process, this structure also induces longer wormlike micelles in final formation as depicted in Fig. 2. In order to give a rough explanation of this behavior, we calculated the mean-square radius of gyration ( $\langle R_g^2 \rangle$ ) of the hydrophilic block in star-dendritic triblock copolymer as well as time evolution of the number of clusters. The  $\langle R_g^2 \rangle$  of the hydrophilic block is defined as

$$\langle R_g^2 \rangle_{\text{hydrophilic}} = \left\langle \frac{1}{N} \sum_{i=1}^N (r_i - r_{\text{cm}})^2 \right\rangle \quad (10)$$

where  $r_i$  and  $r_{\text{cm}}$  denote the position vector of each particle in the hydrophilic block and the center-of-mass for the whole hydrophilic block. Fig. 4 shows the calculated  $\langle R_g^2 \rangle_{\text{hydrophilic}}$  as a function of simulation steps. It is clear that the sizes of hydrophilic domains increase rapidly to the equilibrium values at the beginning of simulation in both models. In addition, the  $\langle R_g^2 \rangle_{\text{hydrophilic}}$  of model II is less than that of model I, which

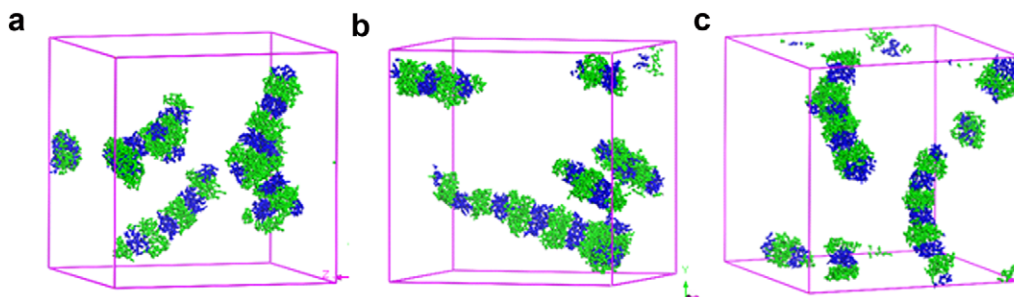


Fig. 2. Morphologies of multicompartment micelles formed from star-dendritic triblock copolymer: (a) model I, (b) model II, and (c) model III. B blocks and water were omitted for clarity; A, blue; C, green (for interpretation of colour in this figure, the reader is referred to the web version of this article).

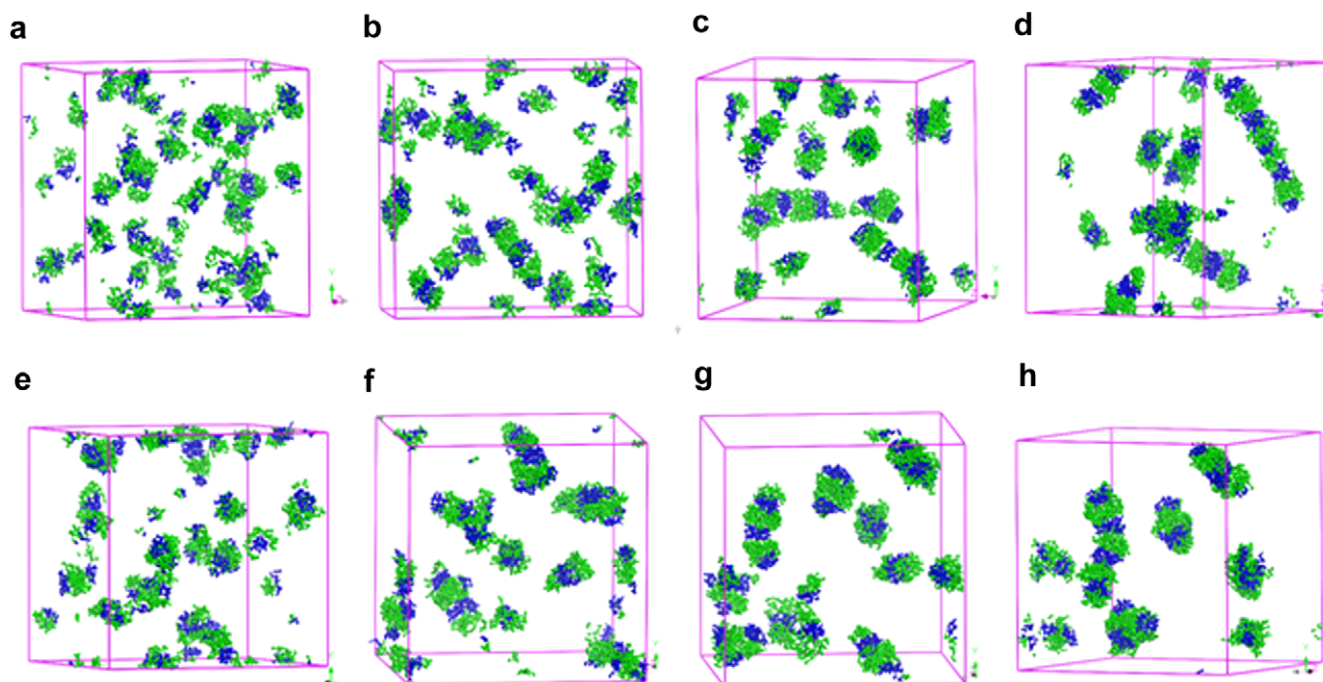


Fig. 3. Evolution of wormlike multicompartiment micelle from different models obtained by DPD simulations: (a)–(d) for model I, (a) step = 5000, (b) step = 10,000, (c) step = 25,000, and (d) step = 35,000; (e)–(h) for model II, (e) step = 10,000, (f) step = 25,000, (g) step = 50,000, and (h) step = 87,000. B blocks and solvents were omitted for clarity; A, blue; C, green (for interpretation of colour in this figure, the reader is referred to the web version of this article).

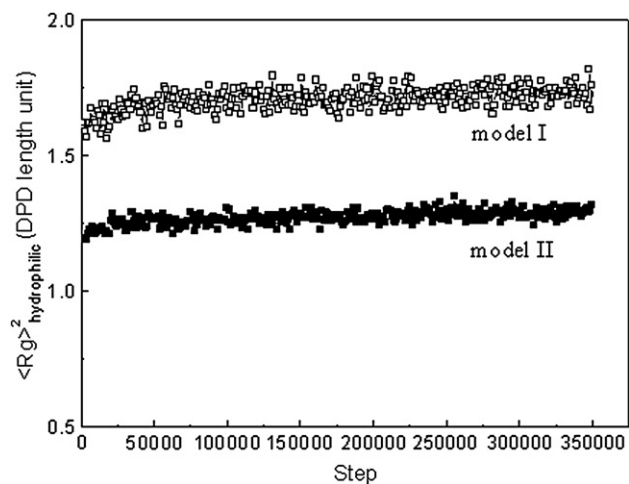


Fig. 4. The mean-square radius of gyration for the two models of star-dendritic triblock copolymers as a function of simulation step.

means that dendritic hydrophilic blocks extend less fully than linear hydrophilic block with the same volume fraction in copolymer. Thus, since there are only a small number of hydrophobic blocks aggregating to form the core of micelle at the beginning of formation process, the compact structure of dendritic hydrophilic blocks is able to prevent the combination of hydrophobic cores and slows down the fusion process. This is corresponded to the fact that model II has larger number of clusters than model I during 1000–125,000 steps as shown in Fig. 5. On the other hand, the hydrophobic effect is commonly considered to be the main driving force in the formation of micelle. The solvents drive the hydrophobic blocks to get together to minimize the contact of hydrophobic blocks and

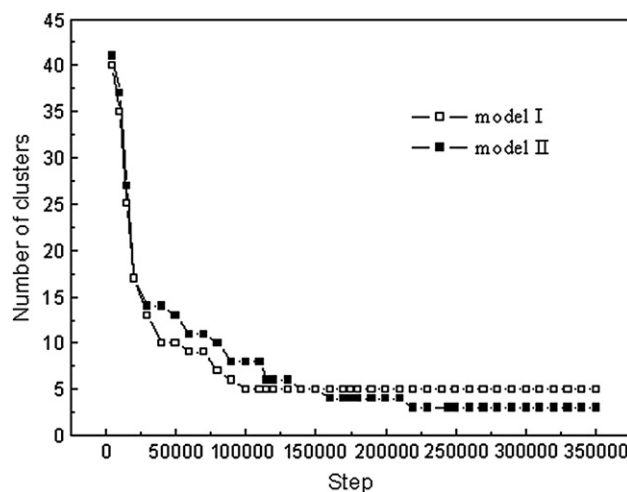


Fig. 5. The number of clusters for the two models of star-dendritic triblock copolymers as a function of simulation step.

solvents, inducing that the core of micelle grows longer. From the previous work, it is concluded that due to the starlike structure of copolymer chains the discrete “hamburger” micelles tend to contact with each other from the top and bottom and they “superpose” one by one to form wormlike micelles with long layered cores [22]. In this case, dendritic hydrophilic block does not extend fully enough to shield the top and bottom of hydrophobic regions from the solvent completely with the core growing. Therefore, the interfacial tension of the micelles goes up and the fusion continues. As a result, the number of clusters in equilibrium state is low in star-dendritic triblock copolymer with dendritic hydrophilic block as shown in Fig. 5 and longer wormlike micelles are formed.



The above simulations show that the dendritic structure in star-dendritic triblock copolymer may affect the detail of final morphology of the multicompartment micelles, and at the same time the formation process of micelle has also been influenced.

### 3.2. Effects of the dendritic structure on the response to the solvent quality of multicompartment micelles

It is well known that one significant application of micelle is delivering drug molecules [1,2], in which the delivering method is a key aspect attracting more and more attention [31,32]. Recently a possible delivering way is proposed based on changing the solvent quality, such as pH value, to destabilize the loading compartment of micelle and then release the drug [33]. Therefore, in this section we attempt to investigate the effects of dendritic structure on the response to the solvent quality, which may provide useful information for the application of multicompartment micelle in the field of drug delivery.

We change the solvent quality by varying the interactions between solvent and hydrophobic block systematically, that is, decreasing the DPD repulsion parameters  $a_{AS}$  from 50 to 30 and  $a_{CS}$  from 120 to 30. It was found that with the decrease of  $a_{AS}$  or  $a_{CS}$  the multicompartment micelle has a morphological transition from compact wormlike structure to elongate wormlike structure, and finally to discrete spherical or “hamburger” micelles, as shown in Fig. 6. In fact, this transition is the gradual dissolution of the corresponding compartments in micelle with decreasing the degree of incompatibility between solvent and hydrophobic block. It is interesting that the morphological change is similar to the reverse course of the fusion process when wormlike multicompartment micelle was formed [22]. The morphologies with different  $a_{AS}$  or  $a_{CS}$  are listed in Table 1. Obviously, the transitions of micelle formed from model II take place at larger values of  $a_{AS}$  or  $a_{CS}$  than that of model I, which means the dendritic structure in hydrophilic block may enhance the stability of the multicompartment micelle when the solvent quality is changed. This is caused by the fact that dendritic structure can form more compact hydrophilic shell outside the multicompartment micelle.

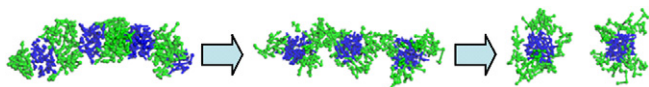


Fig. 6. The snapshots of morphological transition of multicompartment micelle with the change of  $a_{AS}$  or  $a_{CS}$ .

Table 1  
The morphologies of multicompartment micelle with the change of  $a_{AS}$  or  $a_{CS}$

$a_{AS}$	50	40	37	35	33	31
Model I	Compact <sup>a</sup>	Compact	Elongate	Elongate	Discrete	Discrete
Model II	Compact	Compact	Compact	Compact	Elongate	Discrete
$a_{CS}$	120	50	43	40	37	33
Model I	Compact	Compact	Elongate	Elongate	Discrete	Discrete
Model II	Compact	Compact	Compact	Elongate	Elongate	Discrete

<sup>a</sup> Compact, elongate, and discrete represent compact wormlike micelle, elongate wormlike micelle, and discrete micelle, respectively.

To quantitatively understand the difference in structural change of the micelles for models I and II with decreasing the value of  $a_{AS}$  or  $a_{CS}$ , the conservative energy between different parts of system was estimated from the conservative force as follows [34,35]:

$$U_{i<j}^C = \sum_{i<j} \frac{1}{2} a_{ij} (r_c - r_{ij})^2 \quad (11)$$

Since  $a_{ij}$  is same in models I and II, the comparison of conservative energy may give a rough estimation of the difference in distance between particles in different models. We studied two parts of the conservative energy: one is that between block A ( $U_{AA}^C$ ) when  $a_{AS}$  is changed, and the other is that between block C ( $U_{CC}^C$ ) when  $a_{CS}$  is changed. Figs. 7 and 8 show the energy as a function of the value of  $a_{AS}$  or  $a_{CS}$ .

From Fig. 7, we can see that the conservative energy between block A was decreased with the decrease of  $a_{AS}$ . Considering that the conservative force in DPD is soft interaction and only validates within the cutoff radius  $r_c$ , the decrease of energy means that the distance between block A is increased correspondingly and the compartments formed from block A in micelle become elongate, as depicted in the morphological transition in Fig. 6. When  $a_{AS}$  is larger than 40 or less than 33, the conservative energy between block A in model I is similar to that in model II. However, the value of energy in model II is larger than that in model I by about 5% when  $a_{AS}$  is in the range of 33–40. These results indicate that dendritic structure in hydrophilic block has little effect on

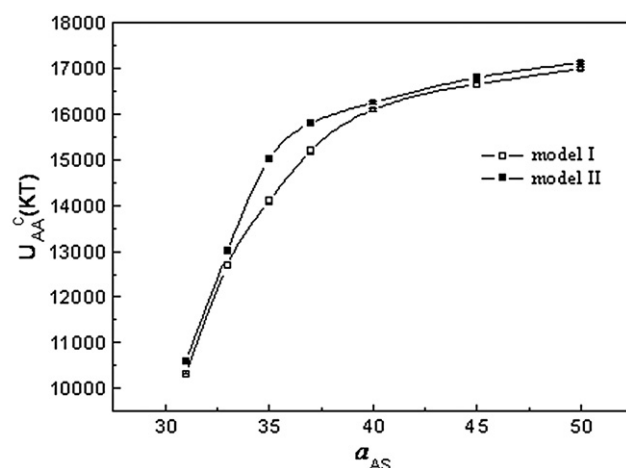


Fig. 7. The conservative energy between block A for the two models as a function of the interaction between block A and solvent.

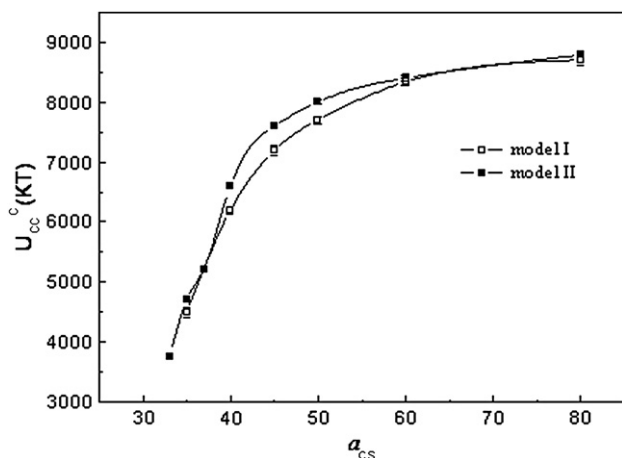


Fig. 8. The conservative energy between block C for the two models as a function of the interaction between block C and solvent.

the structure of hydrophobic core on the condition of compact wormlike micelle or discrete micelle, while this structure may induce more compact core than that of linear block when the corresponding compartment become loose with the change of solvent quality. Similar behavior is also observed when  $a_{CS}$  is varied as shown in Fig. 8.

Furthermore, the mean-square radius of gyration  $\langle R_g^2 \rangle$  as a function of  $a_{AS}$  was calculated as an example to further understand the response to the solvent quality of two models. From the curves in Fig. 9, it can be seen that with the decrease of  $a_{AS}$  from 50 ( $\langle R_g^2 \rangle$  of the copolymer chains has a slight increase, corresponding to the appearance of the elongate wormlike micelles. When the discrete micelles were formed with further decreasing  $a_{AS}$ , the copolymer chains began to contract and  $\langle R_g^2 \rangle$  was decreased. As expected, without shielded by the hydrophilic block with dendritic structure, the increment of  $\langle R_g^2 \rangle$  in model I is larger than that in model II when the elongate wormlike micelles were formed, which is consistent with the difference in the conservative energy between block A as shown in Fig. 7.

The above simulations show that the dendritic structure influences the structural response of multicompartment micelle

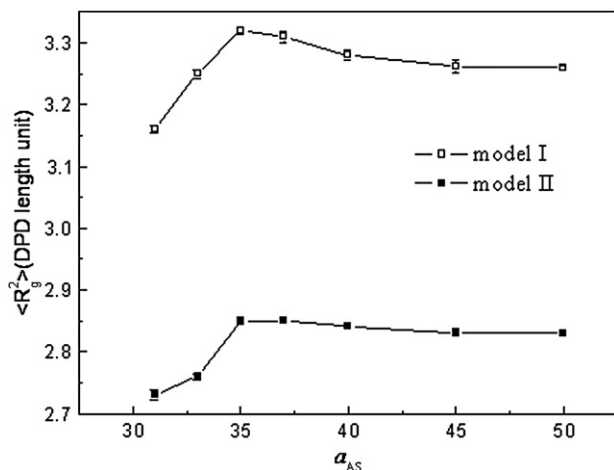


Fig. 9. The mean-square radius of gyration for the two models as a function of the interaction between block A and solvent.

to the change of solvent quality. When the dendritic part is in hydrophilic block, the micelle especially the hydrophobic core is more stable. Thus, the structure of building block copolymer may be chosen on the basis of the quality of target-circumstance where the micelle is used for drug release.

#### 4. Conclusions

In this work, the effect of dendritic structure on the formation as well as the response to the solvent quality of multicompartment micelles was studied in star-dendritic triblock copolymers by DPD method. The results indicate that the dendritic structure not only influences the detailed structure of the hydrophobic core of multicompartment micelles, but also the corresponding formation speed. In addition, dendritic structure in hydrophilic block can enhance the stability of micelle when the solvent quality was changed. Thus, choosing different dendritic block in building copolymers can generate various multicompartment micelles with different response to the solvent quality. This may provide information for the complete understanding of multicompartment micelles, as well as the design and application of this novel kind of micelles in the field of drug delivery.

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